

ALEXANDER LIYVAPUU

Natural vibrations of elastic stepped
arches with cracks



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UNIVERSITY OF TARTU
Press

Institute of Mathematics and Statistics, Faculty of Science and Technology,
University of Tartu, Estonia

Dissertation has been accepted for the commencement of the degree of Doctor of Philosophy (Mathematics) on June 17, 2016 by the Council of the Institute of Mathematics and Statistics, Faculty of Science and Technology, University of Tartu.

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Commencement will take place on August 24, 2016, at 13.00 in J. Liivi 2-403.

Publication of this dissertation has been granted by University of Tartu.

ISSN 1024-4212

ISBN 978-9949-77-175-2 (print)

ISBN 978-9949-77-176-9 (pdf)

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University of Tartu Press
www.tyk.ee

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List of original publications

- I J.Lellep, A.Liyvapuu, Natural vibrations of elastic arches with cracks.
In: Proceedings of the 2nd International Conference Optimization and Analysis of Structures. Editors: J.Lellep, E.Puman, University of Tartu Press, pp. 64–69, 2013.
- II J.Lellep, A.Liyvapuu, Free vibrations of elastic laminated arches.
In: Proceedings of the 3rd International Conference Optimization and Analysis of Structures. Editors: J.Lellep, E.Puman, University of Tartu Press, pp. 52–58, 2015.
- III A.Liyvapuu, J.Lellep, Natural vibrations of stepped arches.
International Journal of Advancements in Mechanical and Aeronautical Engineering 3(1), pp. 68–72, 2016.
- IV J.Lellep, A.Liyvapuu, Natural vibrations of stepped arches with cracks.
Agronomy Research 14(3), pp. 821–830, 2016.
- V J.Lellep, A.Liyvapuu, Free vibrations of uniform and hollow-sectional elastic arches, *Elsevier, Procedia Engineering, Modern Building Materials, Structures and Techniques*, 2016, (accepted).

Author's contribution

The author of this dissertation is responsible for majority of the research in all phases (including writing, simulation and preparing of images) of the papers I–V. The solution procedure was developed in co-operation with supervisors; the statement of the problem belongs to the supervisor.

1 Introduction

1.1 Overview of the literature

Plates, shells and beams (straight and curved) are the main structural elements in more complex structures. Modern construction techniques provide for the production of basic elements with the subsequent installation. Extremely important is the quality control during manufacturing and installation. Here the important role play non-destructive testing methods. One possible method might be offered via an analysis of natural vibrations. Mindlin shows in the [55] that the rotational inertia effect is small for the lower plate modes. Huang [28] studied the effects of rotational inertia and shear deformation for the natural frequencies and mode shapes in uniform Timoshenko beam with boundary conditions for free supported ends. He has shown that the effect of these two effects for the natural frequencies increases with the number of modes or cross-sectional dimensions. But the comparative impact on the normal modes of vibration, seems to be very little. Murty [56] developed a linear approximate equation for the transverse vibrations of uniform short beams, including shear deformation and rotary inertia effects. Its values of natural frequencies are in better agreement with the experimental data for comparison with the results obtained using shift correction coefficients proposed by Timoshenko [72], [70], [71] and Cowper [18]. Adams and Bacon [1] showed that the effect of the shear deformation is a function of the aspect ratio (i.e., ratio of length to thickness) and less than 1% for isotropic materials with an aspect ratio of more than twenty.

The effect of the shear deformation and rotatory inertia on the vibrations of beams, rings and arches are investigated also by Rao [64], Lee *et al.* [40], [42], [41], Wu and Chiang [78], [79], [80]. The work by Shames [68] gives us a reason to use the Euler-Bernoulli beam theory in the dissertation and to assume that plane cross sections, normal to the neutral axis before deformation, continue to remain plane and normal to the neutral axis and do not undergo any strain in their planes. Free vibrations of curved beams or arches in term of Euler-Bernoulli theory are the main subject of research in the dissertation. We will use the first modes of vibrations only in our study. The frequencies of free vibrations of curved beams and arches of constant dimensions of the cross section of the arch are calculated by Petyt and Fleischer [61], Markus and Nanosi [53], Chidaparam and Leissa [15], Cerri *et al.* [12], [13], Auciello and DeRosa [4], DeRosa [20], Viola *et al.* [77], [76], [75]

making use of analytical and numerical methods.

Finite element and finite strip methods have been used by Cheung *et al.* [14], Yang and Chen [84], Yang *et al.* [83], Ishaguddin *et al.* [29], Karaagac *et al.* [32]. A method based on curved beam elements is developed by Wu and Chiang [78], [79], [80].

Due to their high stiffness and strength and other mechanical properties compared to the weight the composite and laminate structures have gained popularity during last decades. This involves the need for investigation of the behavior of structures made of composites and laminates. The foundation of the mechanics of composite materials are presented in the books by Herakovich [27], Jones [30], Daniel [19], Tuttle [73]. Various numerical and semi-analytical models for the analysis of granular materials are developed by Kacianauskas *et al.* [31].

Orthotropic beams and plates are studied by Bui *et al.* [9], Ogierman and Kokot [59], Bao *et al.* [6] and others.

Thin-walled plate and shell structures are treated by Kollar and Springer [36], Reddy [65], Vinson and Sierakovski [74], Qatu [62], Gürdal *et al.* [25], Xiang and Wang [81].

The vibrations and stability of structural elements containing cracks and other defects have deserved the attention of many researchers. However, the most of attention is paid to the vibrations of beams (see Dimarogonas [21], [22], [23], Nandwana and Maiti [58], Nahvi and Jabbari [57], Kisa and Brandon [34], Kisa *et al.* [34], Zheng and Kessissoglou [86]). Lellep and Kägo [43] investigated the influence of defects on the eigenfrequencies of classic stretched strips.

The presence of cracks and other defects in structural elements is a source of additional compliance. Anifantis and Dimarogonas [3], Dimarogonas [22], Chondros *et al.* [16], [17], Rizos *et al.* [66], Kukla [37] explored the idea of an elastic spring modelling the additional flexibility due to a crack in the cases of vibrating beams and bars weakened with crack. This approach was extended to the case of elastic plate strips by Lellep and Kägo [43], [38]. Axisymmetric vibrations of elastic circular cylindrical shells with piece wise constant thickness were treated by Lellep and Roots [49], [50]. In the 1990s scientists have paid attention to the development of methods for describing the vibrations of beams containing a crack. Rizos *et al.* [66], Liang *et al.* [51], [52], Nandwana and Maiti [58], Kisa *et al.* [34], [35], Yang and Chen [84], Ostachowicz and Krawczuk [60], Bamnios and Trochides [5] studied the behavior of a beam having a crack at a certain place. Rizos *et al.* [66],

Ostachowicz and Krawczuk [60], Kisa *et al.* [34] and other attempted to simulate the presence of a crack in a beam using the analogy of rotational springs. According to this approach considered cracked beam is consisting of two parts connected by a rotational spring in place of a crack location. This model has been successful in order to describe the local compliance due to the presence of cracks. In the works Dimarogonas [22], Dimarogonas and Paipetis [23] calculations showed a very good approximation methods of linear elastic fracture mechanics (see Anderson [2], Broberg [7] and Broek [8]) experimental data.

This idea is explored in the non-destructive detection of defects. Another method for non-destructive evaluation of plate-like structures is based on Lamb waves (see Ratassepp *et al.* [63]). Lamb waves are used also in the determination of material constants in [39].

The influence of the axial force on the vibrations of Euler-Bernoulli beams was revealed by Caddemi and Calio [10], also by Matsunaga [54]. The Timoshenko arches are investigated by Calio *et al.* [11]. Here, as well as in the paper by Kawakami *et al.* [33] out-of-plane vibrations of curved beams are considered. Functionally graded materials are considered by Eroglu [24].

The foundations of vibrations of beams, plates and shells are presented in the books by Qatu [62], Reddy [65], Vinson and Sierakowski [74], Henrych [26]. The book by Soedel [69] contains the governing equations of the dynamic analysis of beams, plates, shells and of certain non-shell structures. The main principles of the analysis of composite and laminated structures can be found in the books by Tuttle [73], Qatu [62], Daniel and Ishai [19], Jones [30], Herakovich [27], Sadeghpour *et al.* [67] and others. Arches are considered by Xu *et al.* [82]. In the paper by Lellep and Liyvapuu [45] a method of determination of natural frequencies of elastic laminated arches was developed. These results are extended to the case of composite and laminated arches of variable thickness.

During last decades a lot of attention has been paid to the natural vibrations of beams, plates and shells with defects (see Kukla [37], Yang and Chen [84], Lellep and Roots [49], [50]). In the investigations of this type the key problem is the modelling of influence of the defect on the structural behavior. Dimarogonas [21], also Chondros *et al.* [17] have used the model of "massless rotational spring". According to this concept a beam with a defect or a crack is modelled as a mechanical system consisting of two bars which are connected at the cross section where the defect occurs. This approach was employed by Zheng and Fan [85] and others. Lellep and Liyvapuu [46][47],

[48] extended this method to the case of natural vibrations of circular arches having rectangular cross sections with piece-wise constant dimensions. In the present study the natural frequencies of elastic arches with step wise varying cross sections are investigated.

1.2 Aim of the dissertation

The main goal of the research is to investigate the frequency of natural vibrations of elastic stepped arches with crack-like defects and to analyse the sensitivity of the frequency to the geometrical, mechanical and physical properties of the arch. The reasons to choose the methodology that is used in the study and the detailed description of its development are revealed in the introductory section and in Sections 2, 3 and 4 of the present dissertation.

1.3 Structure of the dissertation

The thesis is organized as follows. The first section (Introduction) provides an overview of the study and analysis of vibration, in particular the analysis of free vibrations, in recent decades; it is followed by the explanation of the aim and structure of the dissertation. The main part of the thesis (Sections 2, 3 and 4) describes in detail our method, which is then applied to the arches having both isotropic and laminated structure. The second section discusses two particular cases of stepped arches. At first, the arch of constant thickness simply supported at both ends is studied. The second case is related to laminated arches of constant thickness without cracks. The cases of simply supported arches and arches clamped at both ends are investigated. The third section addresses the issue of the dependence of the frequency of free vibrations on the parameters of the stepped arches. The fourth section examines the frequency of free vibrations of stepped arches with cracks clamped at both ends. The arches are supposed to be with an internal cavity. Finally, the conclusions in English and Estonian are presented, copies of published articles and the author's CV are enclosed.

2 Free vibrations of arches with cracks

In the second section of the dissertation we consider arches with constant thickness. The determination of natural frequencies of these arches is the goal of the research. At first we consider a simply supported at the both ends circular arch with a crack. Then we will study natural vibrations of a laminated arch simply supported at the both ends and arches clamped at the both ends.

2.1 Elastic arch with a crack

Formulation of the problem

Let us consider natural vibrations of an elastic arch of radius $R = \text{const.}$ It is assumed that the arch is simply supported at both ends (see Fig.1). We use polar coordinates. When using polar coordinates the position of current cross section is defined by the angle φ whereas the edges of the arch correspond to $\alpha = 0$ and $\alpha = \beta$, respectively. It is assumed that at $\varphi = \alpha$ a crack-like defect is located. The defect is treated as a stable crack; no attention will be paid to its extension. Let the depth (length) of the crack be $c < h$. It is assumed herein that the arch has rectangular cross section with dimensions h (thickness) and b (width of the arch). The aim of the section is to determine the frequencies of natural vibrations and to reveal the sensitivity of eigenfrequencies to crack parameters. Note that the arch will be treated as a curved beam with the neutral curved axis lying wholly in one plane. It is assumed that the motion of every point of the neutral curve takes place in this plane only.

Solution of the equation of motion

Equilibrium conditions of an element of the arch lead to the equations (Soedel [69]):

$$\begin{aligned} \frac{\partial M}{\partial s} - Q &= 0, \\ \frac{\partial N}{\partial s} + \frac{Q}{R} + p_s &= \bar{\rho} h \ddot{U}, \\ \frac{\partial Q}{\partial s} - \frac{N}{R} + p_n &= \bar{\rho} h \ddot{W}, \end{aligned} \tag{2.1}$$

where dots denote the differentiation with respect to time t and s is the length of the arch. In (2.1) U and W denote the displacements in the cir-

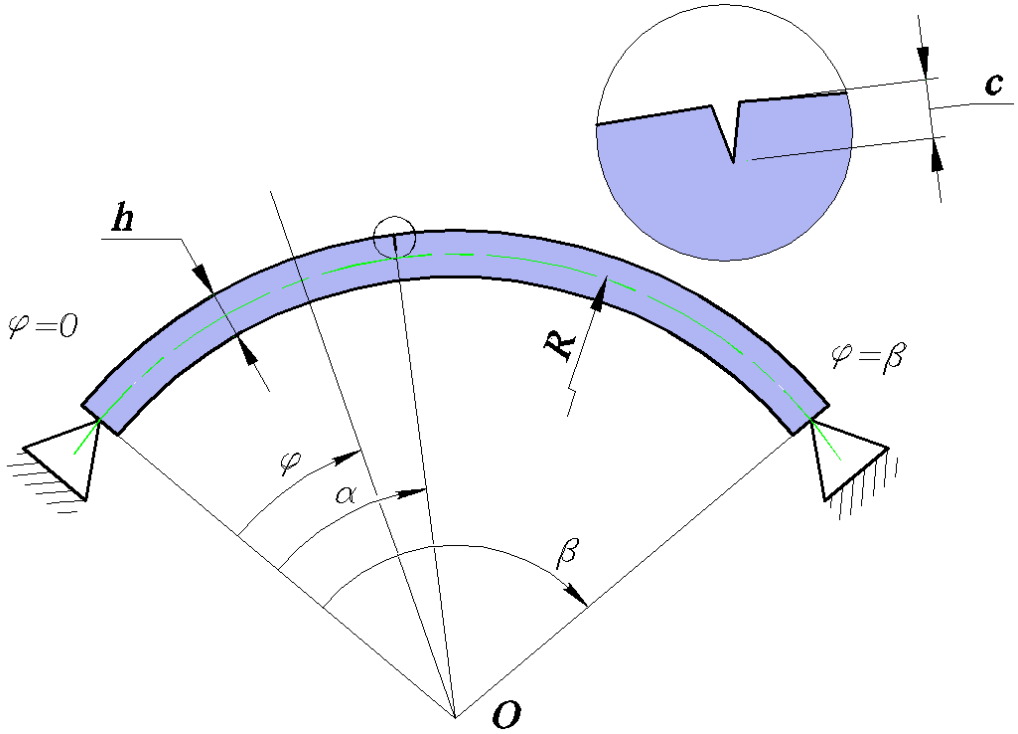


Figure 1: Elastic arch with the crack

cumferential and transverse direction, respectively, whereas p_s and p_n stand for the external loads in these directions. Here N and M are the membrane force and bending moment, Q denotes the shear force; $\bar{\rho}$ is the mass per unit width and h — the thickness of the arch. According to the classical approach $p_s = u = 0$. Assuming that

$$M(0, t) = 0, \quad M(\beta, t) = 0 \quad (2.2)$$

and the membrane force N vanishes at the edges of the arch it follows from (2.1) that $M = -NR$ and

$$M'' + M + R^2(p_n - \bar{\rho}h\ddot{W}) = 0, \quad (2.3)$$

where primes denote the differentiation with respect to the angle φ and $s = \varphi R$.

In the case of curved beams and arches the strain components are the relative extension (see Soedel [69])

$$\varepsilon = \frac{1}{R}(U' + W) \quad (2.4)$$

and the curvature

$$\varkappa = -\frac{1}{R^2}(-U' + W''). \quad (2.5)$$

According to the Hooke's law (see [44], [45], [69])

$$N = Ebh\varepsilon, \quad M = \frac{Eh^3b}{12}\varkappa. \quad (2.6)$$

Here E stands for the Young's modulus.

Thus

$$\begin{aligned} N &= \frac{Ehb}{R}(U' + W), \\ M &= \frac{Eh^3b}{12R^2}(U' - W''), \end{aligned} \quad (2.7)$$

Since $N = 0$ the equation (2.17) yields $U' = -W$. Therefore

$$\varkappa = -\frac{1}{R^2}(W + W'') \quad (2.8)$$

and

$$M = \frac{Eh^3b}{12R^2}(-W - W''). \quad (2.9)$$

We will focus on the free vibrations of the arch. Taking $p_n = 0$ and substituting (2.9) in (2.3) one obtains the equation

$$\frac{Eh^3}{12R^2}(W^{IV} + 2W'' + W) + \bar{\rho}hR^2\ddot{W} = 0. \quad (2.10)$$

In order to solve the differential linear equation (2.10) with partial derivatives the method of Fourier' will be employed. According to this method let us assume that

$$W(\varphi, t) = w(\varphi) \cdot \sin(\omega t), \quad (2.11)$$

where ω stands for the frequency of natural vibrations and $w(\varphi)$ is an unknown function of the variable φ . The substitution (2.11) in (2.10) leads to the equation

$$w^{IV} + 2w'' + w(1 - k^2) = 0, \quad (2.12)$$

where

$$k = \sqrt{\frac{12\bar{\rho}R^4\omega^2}{Eh^2}}. \quad (2.13)$$

The roots of the characteristic equation of the linear differential equation (2.12) are

$$\lambda = \pm\sqrt{-1 \pm k}. \quad (2.14)$$

We assume that $-1 + k > 0$ and $-1 - k < 0$. Thus the general solution of (2.12) takes form

$$w = A_1 \cosh(\mu\varphi) + A_2 \sinh(\mu\varphi) + A_3 \cos(\nu\varphi) + A_4 \sin(\nu\varphi) \quad (2.15)$$

for $\varphi \in [0, \alpha]$ and

$$w = B_1 \cosh(\mu\varphi) + B_2 \sinh(\mu\varphi) + B_3 \cos(\nu\varphi) + B_4 \sin(\nu\varphi) \quad (2.16)$$

for $\varphi \in [\alpha, \beta]$.

Here the notation $\mu^2 = 1 - k$, $\nu^2 = 1 + k$ is used.

Arbitrary constants A_1, A_2, A_3, A_4 and B_1, B_2, B_3, B_4 in (2.15), (2.16) have to be determined from boundary and intermediate conditions.

Additional flexibility caused by the crack

Let us use the method developed by Chondros *et al.* [17], Dimarogonas [22], Kukla [37]. The authors of [17] supposed that a slope of the deflection can be considered as a discontinuous quantity at the cross section with cracks.

Let us denote the jump of the discontinuous quantity at the cross section $\varphi = \alpha$ as

$$\theta = W'(\alpha + 0, t) - W'(\alpha - 0, t). \quad (2.17)$$

It is recognized that there exists a relationship between the local compliance of the arch C and the stress intensity factor K known in the linear elastic fracture mechanics.

The quantity θ can be treated as a generalized displacement corresponding to the generalized force (bending moment) $M_c = M(\varphi, t)|_{\varphi=\alpha}$ one can write

$$\theta = C \cdot M_c, \quad (2.18)$$

where C is the additional compliance caused by the crack.

On the other hand, it is known in the fracture mechanics that the energy release rate can be calculated as (here $A = bc$)

$$G = \frac{M_c^2}{2} \frac{dC}{dA} \quad (2.19)$$

or as (see Anderson [2], Broek [8])

$$G = \frac{K^2}{E'}. \quad (2.20)$$

In (2.20) $E' = E$ for plane stress state and $E' = \frac{E}{(1-\nu^2)}$ for the plane deformation state.

The stress intensity factor K itself is defined as

$$K = \sigma \sqrt{\pi c} F(s). \quad (2.21)$$

Here $s = \frac{c}{h}$. If the cracked element is subjected to pure bending then

$$\sigma = \frac{6M_c}{bh^2}. \quad (2.22)$$

The shape function F in (2.21) must be defined on the basis of experimental results. It depends on the type of a cracked element. Following Dimarogonas [21], Rizos *et al.* [66] one can take

$$F(s) = 1.93 - 3.07s + 14.53s^2 - 25.11s^3 + 25.80s^4, \quad (2.23)$$

where $s = c/h$.

Combining the relations (2.18)—(2.23) after algebraic manipulations one obtains

$$\frac{dC}{ds} = \frac{72\pi}{E'h^2b} s F^2(s). \quad (2.24)$$

The integration of (2.24) leads to the result

$$C = \frac{72\pi}{E'h^2b} f(s), \quad (2.25)$$

where

$$f(s) = \int_0^s \xi F^2(\xi) d\xi. \quad (2.26)$$

Thus the jump of the slope can be evaluated as

$$w'(\alpha + 0) - w'(\alpha - 0) = p \left(w''(\alpha) + w(\alpha) \right), \quad (2.27)$$

where

$$p = -\frac{6\pi h(1 - \nu^2)}{R^2} \cdot f(s). \quad (2.28)$$

Boundary and intermediate conditions for a simply supported arch

Boundary conditions for an arch simply supported at both ends are

$$\begin{aligned} w(0) &= 0, & w''(0) &= 0, \\ w(\beta) &= 0, & w''(\beta) &= 0. \end{aligned} \quad (2.29)$$

At the internal points of the interval $(0, \beta)$ the displacement W and its derivative, also bending moment M and shear force Q must be continuous. Therefore one has

$$\begin{aligned} w(\alpha - 0) &= w(\alpha + 0), \\ w''(\alpha - 0) &= w''(\alpha + 0), \\ w'''(\alpha - 0) &= w'''(\alpha + 0). \end{aligned} \quad (2.30)$$

The quantity w' is discontinuous at $\varphi = \alpha$; the jump condition for it (2.27) must be fulfilled.

The conditions (2.29) furnish at $\varphi = 0$ the equations

$$\begin{aligned} A_1 + A_3 &= 0, \\ \mu^2 A_1 - \nu^2 A_3 &= 0. \end{aligned} \quad (2.31)$$

It immediately follows from (2.31) that

$$A_1 = A_3 = 0. \quad (2.32)$$

The conditions (2.29) furnish at $\varphi = \beta$ the equations

$$B_1 = -B_2 \tanh(\mu\beta), \quad B_3 = -B_4 \tan(\nu\beta). \quad (2.33)$$

Taking (2.27), (2.30), (2.32), (2.33) into account the continuity requirements can be presented in the form of the following system

$$\begin{aligned}
A_2 \sinh(\mu\alpha) + A_4 \sin(\nu\alpha) &= \\
&= B_2 \left(\sinh(\mu\alpha) - \tanh(\mu\beta) \cosh(\mu\alpha) \right) + \\
&\quad + B_4 \left(\sin(\nu\alpha) - \tan(\nu\beta) \cos(\nu\alpha) \right), \\
A_2 \mu \cosh(\mu\alpha) + A_4 \nu \cos(\nu\alpha) &= \tag{2.34} \\
&= p \left[B_2(\mu^2 + 1) \left(\sinh(\mu\alpha) - \tanh(\mu\beta) \cosh(\mu\alpha) \right) + \right. \\
&\quad \left. + B_4(1 - \nu^2) \left(\sin(\nu\alpha) - \tan(\nu\beta) \cos(\nu\alpha) \right) \right],
\end{aligned}$$

$$\begin{aligned}
A_2 \mu^2 \sinh(\mu\alpha) - A_4 \nu^2 \sin(\nu\alpha) &= \\
&= B_2 \mu^2 \left(\sinh(\mu\alpha) - \tanh(\mu\beta) \cosh(\mu\alpha) \right) + \\
&\quad + B_4 \nu^2 \left(\tan(\nu\beta) \cos(\nu\alpha) - \sin(\nu\alpha) \right),
\end{aligned}$$

$$\begin{aligned}
A_2 \mu^3 \cosh(\mu\alpha) - A_4 \nu^3 \cos(\nu\alpha) &= \\
&= B_2 \mu^3 \left(\cosh(\mu\alpha) - \tanh(\mu\beta) \sinh(\mu\alpha) \right) + \\
&\quad + B_4 \nu^3 \left(-\cos(\nu\alpha) - \tan(\nu\beta) \sin(\nu\alpha) \right).
\end{aligned}$$

Evidently, we have the linear homogeneous system of four algebraic equations (2.34) with the unknowns A_2, A_4, B_2, B_4 . This system has a non-trivial solution if and only if its determinant Δ equals to zero.

Numerical results and discussion

The equation $\Delta = 0$ is solved up to the end numerically making use of the computer code MATLAB. We found the lowest frequency corresponding to the first mode of vibration. Some examples are shown in Fig.2, Fig.3. The results presented in Fig.2, Fig.3 correspond to the simply supported arch made of steel with moduli $E = 2.1 \cdot 10^{11} \text{ Pa}$, $\nu = 0.3$, $\rho = 7865 \text{ kg/m}^3$.

The geometrical dimensions of the arch are $R = 1 \text{ m}$, $b = 0.05 \text{ m}$, $h = 0.05 \text{ m}$. In Fig.2 $\beta = 1.5 \text{ rad}$ and in Fig.3 $\beta = 1 \text{ rad}$.

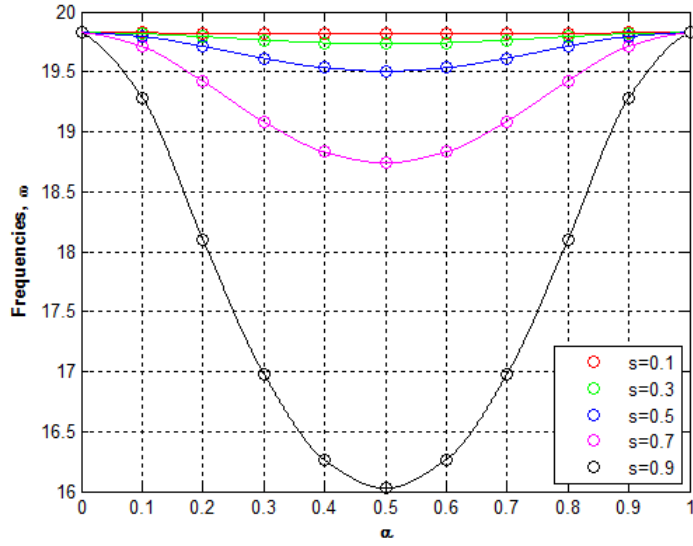


Figure 2: Natural frequency vs. length of the crack, $\beta = 1.5$ rad.

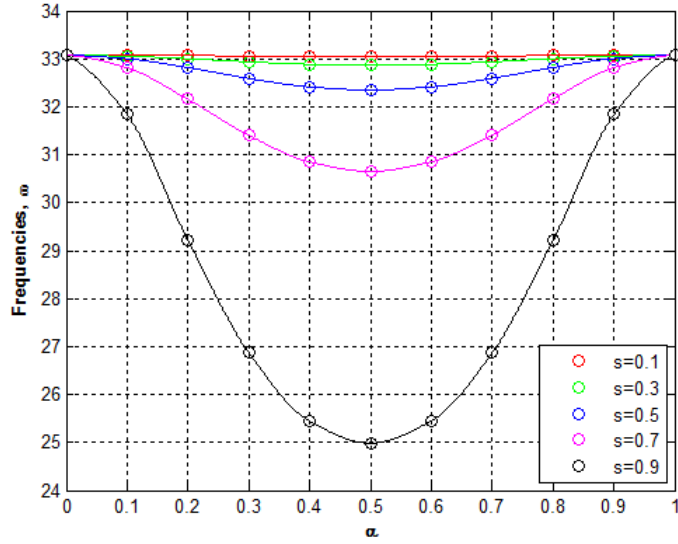


Figure 3: Natural frequency vs. length of the crack, $\beta = 1$ rad.

Different curves in Fig.2 and Fig.3 correspond to different extensions of the crack. Evidently, the higher frequencies correspond to the larger angle of the arch. It can be seen from Fig.2 and Fig.3 that the eigenfrequency has the highest level for the arch without any defects.

Results obtained by the method presented above are compared with the results of Wu, Chang [78] and Petyt, Fleischer [61] in the case of simply supported arches in Table 2.1. and Table 2.2.

Here K indicates the number of natural frequency.

Table 2.1 corresponds to the arch with $\beta = 1 \text{ rad}$, $R = 30 \text{ in} \approx 0.762 \text{ m}$, $h = 0.01289 \text{ in} \approx 3.27 \cdot 10^{-4} \text{ m}$

The material of the arch is a mild steel with characteristics:

$E = 10^7 \text{ lb/in}^2 \approx 2.871 \cdot 10^{11} \text{ Pa}$, $\rho = 0.1 \text{ lb/in}^3 \approx 2768.2 \text{ kg/m}^3$.

Here $\omega_0 = \omega \sqrt{\frac{AR^4}{EI}}$, $I = \frac{1}{12}bh^3$, $A = bh$.

Table 2.1. Comparison of results, ($\beta = 1 \text{ rad}$)

K	1	2	3	4	5
ω_0	0.366	1.590	3.672	6.481	10.159
Petyt, Fleischer [61]	0.349	1.571	3.612	6.470	10.144

Table 2.2. Comparison of results, ($\beta = \pi/2$)

K	1	2	3	4	5
ω_0	2	15	35	63	99
Wu, Chiang [78]	—	13.773	32.426	61.610	96.375

It can be seen from Table 2.1 and Table 2.2 that the results of the current study are quite close to results obtained by other researchers.

2.2 Laminated arches

Basic assumptions

Free vibrations of an elastic arch made of a multi layered material will be considered. It assumed that the cross-section of the arch is rectangular with the width b and total height (thickness) $H = \text{const}$.

The cross section of the arch consists of layers with thickness h_j ($j = 0, 1, \dots, n$). Each layer is assumed to be an elastic layer with material pa-

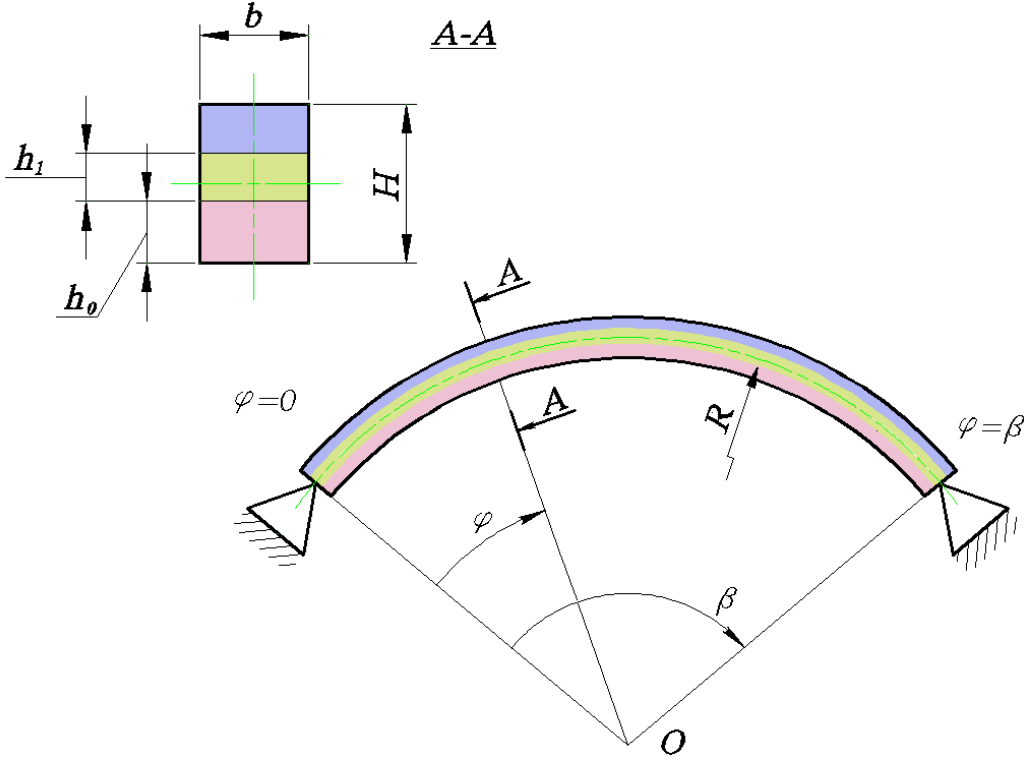


Figure 4: Elastic laminated arch

rameters ρ_j (density), E_j (Young's modulus), ν_j (Poisson's ratio). The layers are treated as orthotropic laminae consisting of a matrix material and of uniformly embedded fibers. However, in the case of an orthotropic lamina these engineering constants are not independent. For a plane stress situation one has (see Gürdal *et al.*[25]) for the neighbouring layers must be satisfied restrictions

$$\frac{\nu_{ij}}{E_i} = \frac{\nu_{ji}}{E_j} \quad (2.35)$$

for $i = 0, 1, \dots, n-1$, $j = i+1$.

The goal of the study is to determine the natural frequencies of free vibrations of laminated elastic arches and to elucidate the sensitivity of eigenfrequencies on the geometrical parameters of the laminate.

Governing equations

In the theory of laminated elastic plates the physical relations can be presented as (see Kollar and Springer [36], Reddy [65], Vinson and Sierakowski [74])

$$N = A\varepsilon + B\kappa \quad (2.36)$$

and

$$M = B\varepsilon + D\kappa. \quad (2.37)$$

Here N and M stand for vectors of membrane forces and principal moments, respectively, whereas ε and κ are corresponding vectors of strain components and curvatures. The elements of matrices A , B and D can be calculated as (Vinson and Sierakowski [74])

$$\begin{aligned} A_{ij} &= \sum_{k=1}^n Q_{ij}^{(k)} (z_k - z_{k-1}), \\ B_{ij} &= \frac{1}{2} \sum_{k=1}^n Q_{ij}^{(k)} (z_k^2 - z_{k-1}^2), \\ D_{ij} &= \frac{1}{3} \sum_{k=1}^n Q_{ij}^{(k)} (z_k^3 - z_{k-1}^3), \end{aligned} \quad (2.38)$$

where $|z_k - z_{k-1}| = h_k$.

In the case of laminates for which each lamina is reinforced with a unidirectional array of fibers one has

$$Q_{11}^{(k)} = \frac{E_1^{(k)}}{1 - \nu_{12}^{(k)} \nu_{21}^{(k)}}, \quad Q_{12}^{(k)} = \frac{\nu_{12}^{(k)} E_2^{(k)}}{1 - \nu_{12}^{(k)} \nu_{21}^{(k)}}, \quad Q_{22}^{(k)} = \frac{E_2^{(k)}}{1 - \nu_{12}^{(k)} \nu_{21}^{(k)}}, \quad (2.39)$$

where the superscript (k) indicates the number of the lamina. In the case of beams and arches the quantities A , B and C are scalars. Also the vectors ε and κ are one-dimensional strain components.

These can be expressed as (2.6), (2.5).

The Hooke's law furnishes with (2.36), (2.37) the relations

$$N = \frac{A}{R} (U' + W) + \frac{B}{R^2} (U' - W'') \quad (2.40)$$

and

$$M = \frac{B}{R} (U' + W) + \frac{D}{R^2} (U' - W''). \quad (2.41)$$

According to the Euler-Bernoulli beam theory there is no extension of the middle surface of the arch. Therefore, one can take $\varepsilon = 0$ and $U' = -W$.

Thus one obtains

$$M = D\kappa, \quad (2.42)$$

where

$$D = \frac{Eh^3b}{12}. \quad (2.43)$$

Evidently, (2.42) coincides with (2.9). Similarly, the the equilibrium conditions lead to the equation (2.10), as in the previous case.

Let the general solution on the equation of motion be

$$w = C_1 \cosh(\mu\varphi) + C_2 \sinh(\mu\varphi) + C_3 \cos(\nu\varphi) + C_4 \sin(\nu\varphi), \quad (2.44)$$

where C_1, C_2, C_3 and C_4 are arbitrary constants and

$$\mu = \sqrt{1 - \omega R^2 \sqrt{\frac{\bar{\rho}}{D}}}, \quad \nu = \sqrt{1 + \omega R^2 \sqrt{\frac{\bar{\rho}}{D}}}. \quad (2.45)$$

In the case of an arch simply supported at both ends the boundary requirements can be presented by (2.29).

It follows from (2.29), (2.44), (2.45) that

$$\omega = \pm \sqrt{\frac{D}{\bar{\rho}}} \cdot \frac{\pi^2 k^2 - \beta^2}{\beta^2 R^2}, \quad (2.46)$$

where $k = 1, 2, \dots, n$.

Arches clamped at the both ends

In the case of an arch clamped at both ends the boundary requirements are

$$w(0) = 0, \quad w'(0) = 0 \quad (2.47)$$

and

$$w(\beta) = 0, \quad w'(\beta) = 0. \quad (2.48)$$

Conditions (2.47), (2.48) with (2.44) lead to the linear algebraic system with the determinant

$$\Delta = \begin{vmatrix} 1 & 0 & 1 & 0 \\ 0 & \mu & 0 & \nu \\ \cosh(\mu\beta) & \sinh(\mu\beta) & \cos(\nu\beta) & \sin(\nu\beta) \\ \mu \sinh(\mu\beta) & \mu \cosh(\mu\beta) & -\nu \sin(\nu\beta) & \nu \cos(\nu\beta) \end{vmatrix}. \quad (2.49)$$

Evidently, the determinant (2.49) equals to zero.

Numerical results

The results of calculations are presented in Fig.5—Fig.7 and Table 2.3. The obtained results correspond to the three-layers arch with thicknesses h_0, h_1, h_2 .

The dependence of the natural frequency on the order of stacking of layers with different materials is shown in Fig.5. Here as in the previous case $h_0 = h_1 = h_2 = 0.01 \text{ m} = \text{const}$.

The natural frequency versus radius of the arch is depicted in Fig.6 for different values of the thickness h_2 . Here $h_0 = h_1 = 0.01 \text{ m}$.

It can be seen from Fig.6 that the natural frequency decreases with increasing the radius of the arch.

The sensitivity of the natural frequency on the angle β is shown in Fig.7 for different values of the radius of the arch. Here $h_0 = h_1 = h_2 = 0.01 \text{ m}$. It can be seen from Fig.7 that the larger is the radius of the arch the lower is the natural frequency.

The frequency decreases if the angle β increases.

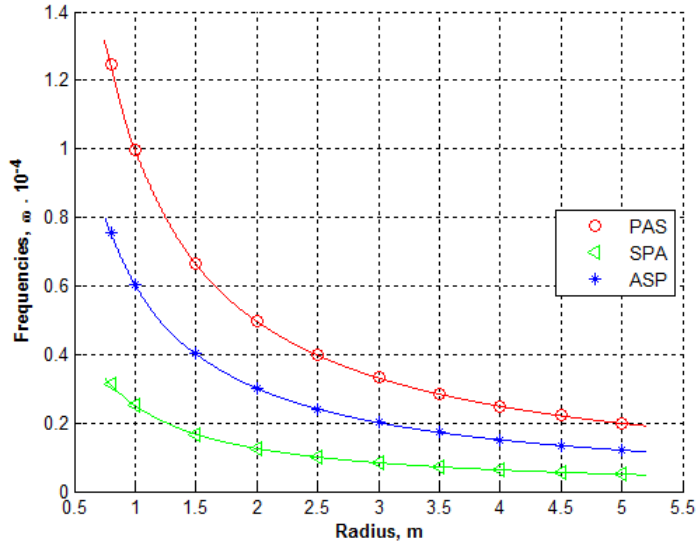


Figure 5: Natural frequency vs. radius of the arch

(A– aluminum, P– polystirol, S– steel)

Table 2.3. Materials of layers

		E, Pa	$\rho, kg/m^3$	ν
(h_2)	Steel	$2.1 \cdot 10^{11}$	7865	0.3
(h_1)	Aluminum	$7 \cdot 10^9$	2700	0.35
(h_0)	Polystirol	$1.5 \cdot 10^9$	30	0.1

It can be seen from Fig.5 that the lowest natural frequency is achieved in the case of stacking materials in the order Steel—Polystirol—Aluminum.

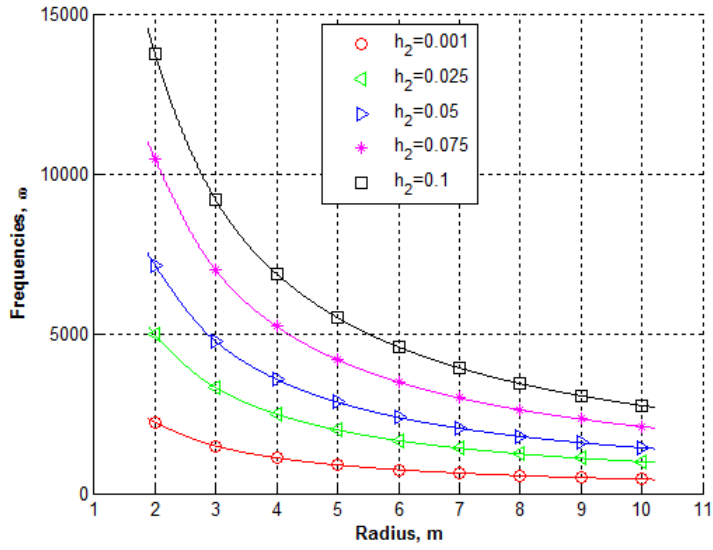


Figure 6: Natural frequency vs. radius of the arch

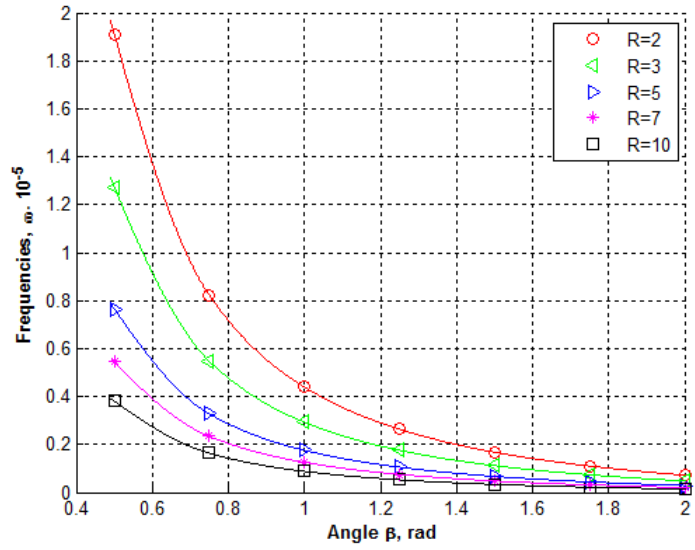


Figure 7: Natural frequency vs. angle of the arch

3 Natural frequencies of stepped arches

The aim of this section is to determine the eigenfrequencies of elastic stepped arches and to study the sensitivity of free vibrations on the crack location and depth. This section is based on the papers Lellep and Liyvapuu [46], [47].

3.1 Stepped arches without defects

Let us consider an elastic circular arch of radius R . The arch is simply supported at the both ends (see Fig.8). As in the previous sections the current angle φ defines the position of the cross-section of the arch. It is assumed that the arch has rectangular cross-section with a constant width b and the height $h = h_j$ for $\varphi \in (\alpha_j, \alpha_{j+1})$. The quantities h_j are assumed to be given constants.

It is assumed that due to the initial excitation the arch is performing free vibrations around its equilibrium position. The goal is to determine the natural frequencies of the free vibrations and to reveal the sensitivity of parameters. We can assume that the arch under consideration is made of a composite of laminated material.

In the theory of elastic plates (see Kollar and Springer [36], Reddy[65], Vinson and Sierakowski[74]) membrane force N and bending moment M are presented as follows

$$N = \sum_{k=1}^{K_j} \int_{z_k}^{z_{k+1}} \sigma_s dz \quad (3.1)$$

and

$$M = \sum_{k=1}^{K_j} \int_{z_k}^{z_{k+1}} \sigma_s z dz,$$

$$|z_{k+1} - z_k| = h_{kj}$$

is the thickness of layer number k ; here z_k is the local coordinate directed through the thickness of the arch and K_j is the number of layers.

Shear force

$$Q_s = \int_{-h/2}^{h/2} \sigma_{s\varphi} dz, \quad (3.2)$$

where $\sigma_{s\varphi}$ — tangential stress component

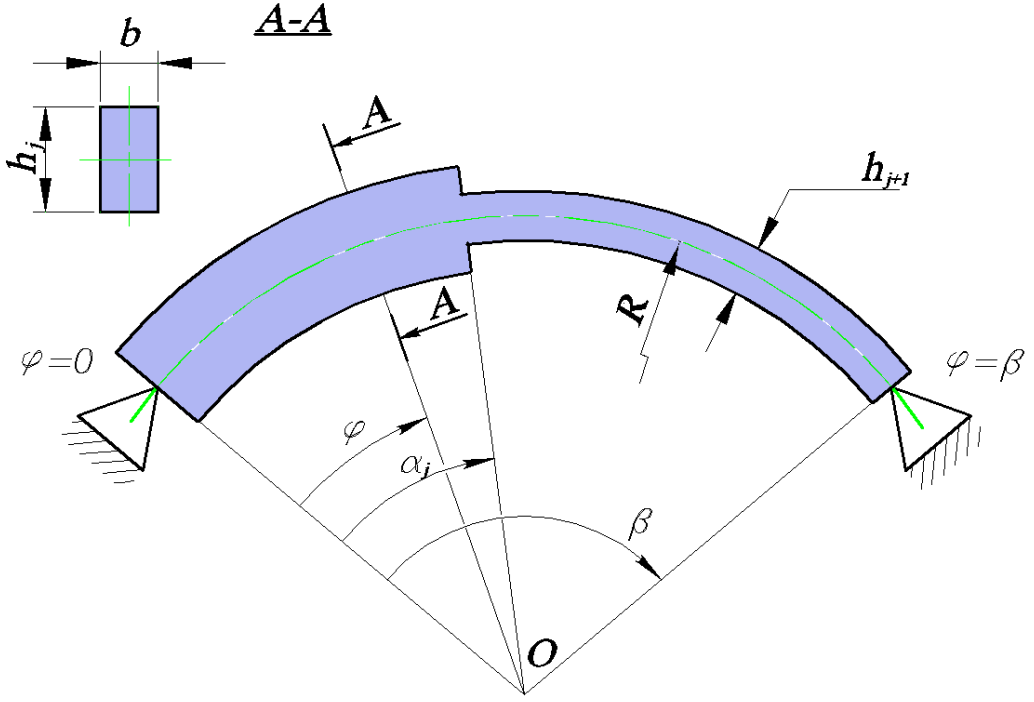


Figure 8: Elastic stepped arch

The Hooke's law for a multilayered arch is given by (2.36), (2.37). Here ε and \varkappa are vectors of strain components and curvatures, respectively, expressed by (2.6), (2.5).

The elements of matrices A, B and D (see Vinson and Sierakowski [74]) given by (2.38).

Also the vectors ε and \varkappa are one-dimensional strain components.

It easily follows from governing equations that in the case of stepped arches and

$$M_j = -D_j \varkappa, \quad (3.3)$$

where \varkappa is defined by (2.8). In the case of free vibrations of stepped arches one has

$$\frac{D_j}{R^2} (W^{IV} + 2W'' + W) + \bar{\rho}_j R^2 \ddot{W} = 0, \quad (3.4)$$

where $\bar{\rho}_j$ — mass per unit area of the middle surface of the arch in different regions.

Here $j = 0$ for $\varphi \in (0, \alpha)$ and $j = 1$ for $\varphi \in (\alpha, \beta)$.

The general solution of equation (3.4) is

$$w = C_{j1} \cosh(\mu_j \varphi) + C_{j2} \sinh(\mu_j \varphi) + C_{j3} \cos(\nu_j \varphi) + C_{j4} \sin(\nu_j \varphi), \quad (3.5)$$

where C_{j1}, C_{j2}, C_{j3} and C_{j4} are arbitrary constants and

$$\mu_j = \sqrt{1 - \omega R^2 \sqrt{\frac{\bar{\rho}_j}{D_j}}}, \quad \nu_j = \sqrt{1 + \omega R^2 \sqrt{\frac{\bar{\rho}_j}{D_j}}}. \quad (3.6)$$

Boundary conditions for an arch simply supported at the both edges are presented by (2.29) and for clamped arches by (2.47), (2.48).

Thus one can write

$$\begin{aligned} w(\alpha_j - 0) &= w(\alpha_j + 0), \\ w'(\alpha_j - 0) &= w'(\alpha_j + 0), \\ D_{j-1} \left(w(\alpha_j - 0) + w''(\alpha_j - 0) \right) &= D_j \left(w(\alpha_j + 0) + w''(\alpha_j + 0) \right), \\ D_{j-1} \left(w'(\alpha_j - 0) + w'''(\alpha_j - 0) \right) &= D_j \left(w'(\alpha_j + 0) + w'''(\alpha_j + 0) \right) \end{aligned} \quad (3.7)$$

for each $j = 1, 2, \dots, n$,

where

$$D_j = \frac{E h_j^3 b}{12 R^2}. \quad (3.8)$$

In order to determine the frequency of the free vibrations of stepped arch one has to determine the constants $C_{j1}, C_{j2}, C_{j3}, C_{j4}$ so that all boundary and intermediate conditions are satisfied. Thus the system of boundary conditions with intermediate conditions (3.7) consists of $4n + 4$ algebraic equations.

In the case $n = 1$ the determinant of this system can be expressed as

$$\Delta = \begin{vmatrix} \xi_1 & \xi_2 & \xi_3 & \xi_4 \\ \mu_0 \xi_5 & \nu_0 \xi_6 & \mu_1 \xi_7 & -\nu_1 \xi_8 \\ \mu_0^2 \xi_1 & -\nu_0^2 \xi_2 & \mu_1^2 \xi_3 & -\nu_1^2 \xi_4 \\ \mu_0^3 \xi_5 & -\nu_0^3 \xi_6 & \mu_1^3 \xi_7 & \nu_1^3 \xi_8 \end{vmatrix}, \quad (3.9)$$

where

$$\begin{aligned} \xi_1 &= \sinh(\mu_0 \alpha), & \xi_5 &= \cosh(\mu_0 \alpha), \\ \xi_2 &= \sin(\nu_0 \alpha), & \xi_6 &= \cos(\nu_0 \alpha), \\ \xi_3 &= \tanh(\mu_1 \beta) \cosh(\mu_1 \alpha) - \sinh(\mu_1 \alpha), & \xi_7 &= \tan(\nu_1 \beta) \sin(\nu_1 \alpha) + \cos(\nu_1 \alpha), \\ \xi_4 &= \tan(\nu_1 \beta) \cos(\nu_1 \alpha) - \sin(\nu_1 \alpha), & \xi_8 &= \tanh(\mu_1 \beta) \sinh(\mu_1 \alpha) - \cosh(\mu_1 \alpha). \end{aligned}$$

The equation $\Delta = 0$ is solved numerically using the computer code of MATLAB. The results of calculations are presented in Fig.9—Fig.11.

The lowest frequency of free vibrations is presented in Fig.9 for various values of the angle β . The arch is made of a mild steel with $E = 2.1 \cdot 10^{11} \text{ Pa}$. The dimensions of the arch are $R = 1 \text{ m}$, $h_0 = 0.02 \text{ m}$, $h_1 = 0.01 \text{ m}$.

The natural frequency versus the step coordinate α is presented in Fig.10. Different curves in Fig.10 correspond to different values of the ratio of thicknesses $\gamma = h_1/h_0$.

Here $\beta = 1.5 \text{ rad}$, $b = 0.01 \text{ m}$, $h_0 = 0.05 \text{ m}$.

The relationship between ω and α is shown in Fig.11 for different materials. Here $R = 1 \text{ m}$, $h_0 = 0.02 \text{ m}$, $h_1 = 0.01 \text{ m}$. It can be seen from Fig.11 that curves corresponding to bronze and concrete are quite near to each other.

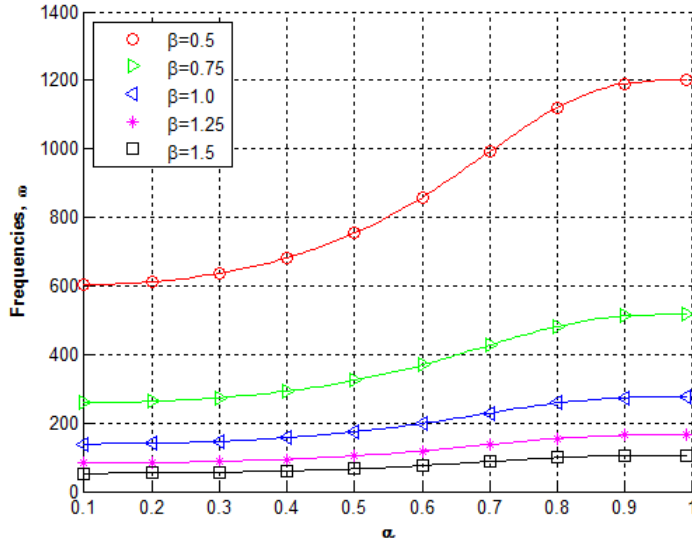


Figure 9: Natural frequency vs. angle of the arch

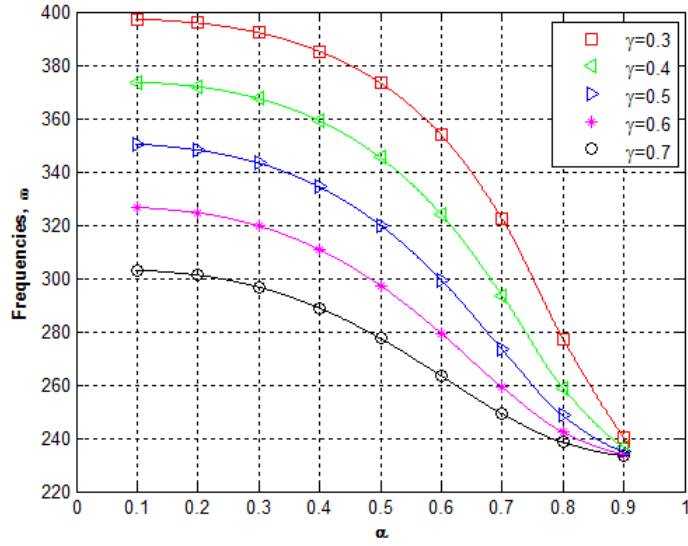


Figure 10: Natural frequency vs. ratio of the thicknesses, $\gamma = \frac{h_1}{h_0}$

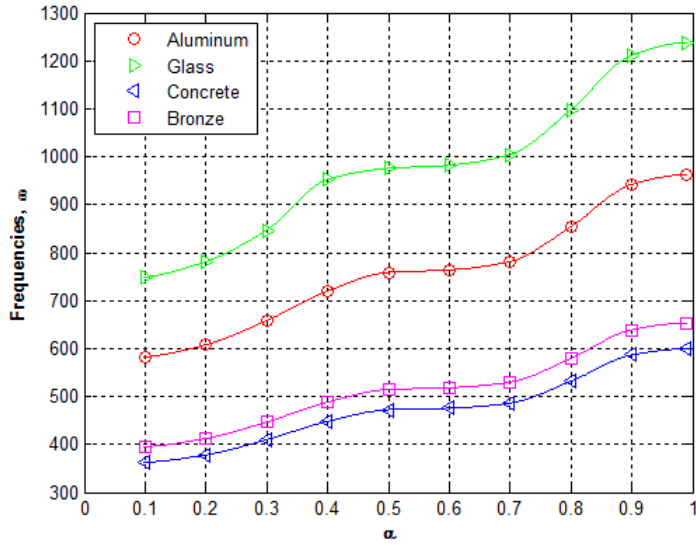


Figure 11: Natural frequency vs. materials of the arch made

3.2 Stepped arch with cracks

Let us consider an elastic stepped arch with cracks. We assume that the arch has piece-wise constant thickness

$$h = h_j, \quad \varphi \in (\alpha_j, \alpha_{j+1}) \quad (3.10)$$

for $j = 0, 1, \dots, n$. Cracks with length c_j are placed in re-entrant corners of steps (see Fig.12). Let other geometrical, physical and mechanical properties of the arch be the same as in the previous section 3.1.

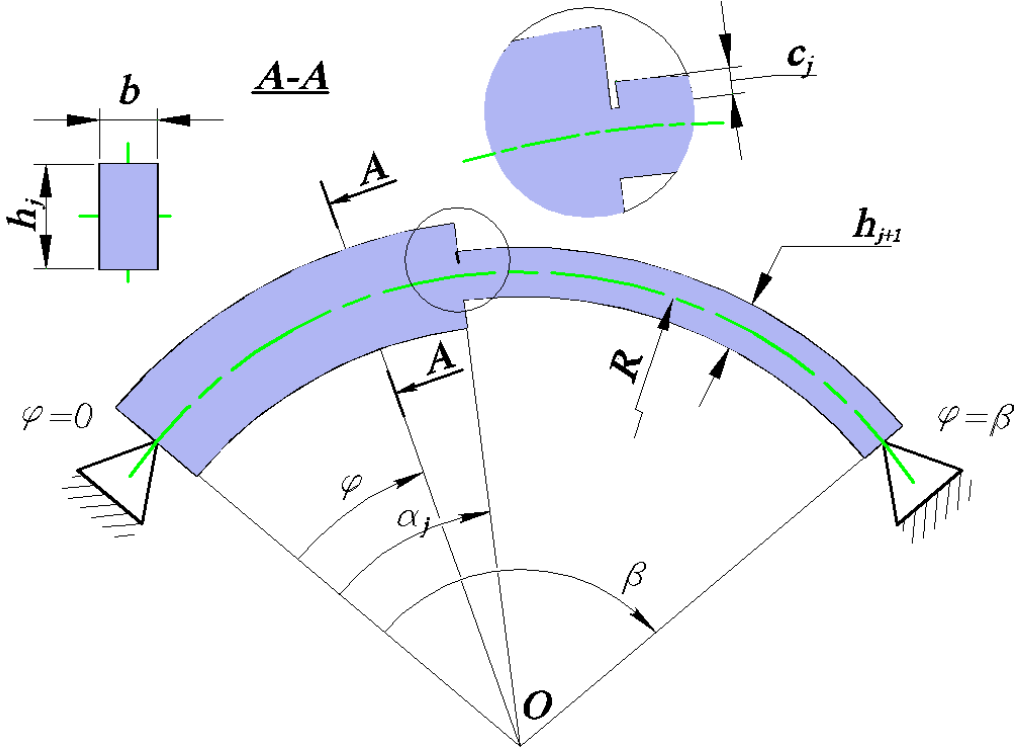


Figure 12: Elastic stepped arch with cracks

Let us study an impact of a surface crack in the cross section $\varphi = \alpha_j$ on the behavior of the arch. Evidently, the continuity conditions and boundary requirements used in the previous section remain valid.

The goal of this section is to study the behavior of an elastic stepped arch with surface cracks.

Equations of equilibrium of an element of a vibrating arch can be transformed into (3.4) for $\varphi \in (\alpha_j, \alpha_{j+1}) \quad j = 0, 1, \dots, n$.

Now the general solution of (3.4) is to be taken as

$$w = C_{1j} \cosh(\mu_j \varphi) + C_{2j} \sinh(\mu_j \varphi) + C_{3j} \cos(\nu_j \varphi) + C_{4j} \sin(\nu_j \varphi). \quad (3.11)$$

For determination of constants C_1, C_2, C_3 and C_4 one can use the boundary conditions (2.29), also intermediate conditions (2.30) and (2.27). The set of intermediate conditions can be presented as (see Leliep and Liyvapuu [44]) in form

$$\begin{aligned} [w(\alpha_j)] &= 0, \\ [w'(\alpha_j)] &= \frac{6\pi h_{j+1}}{R^2} \cdot f(s_j) \left(w''(\alpha_j + 0) + w(\alpha) \right), \\ [w''(\alpha_j)] &= 0, \\ [w'''(\alpha_j)] &= 0. \end{aligned} \quad (3.12)$$

In (3.12) the square brackets denote the jump of corresponding quantities and

$$\begin{aligned} f(s_j) &= 1.86s_j^2 - 3.95s_j^3 + 16.37s_j^4 - 34.23s_j^5 + \\ &+ 76.81s_j^6 - 126.9s_j^7 + 172s_j^8 - 143.97s_j^9 + 66.56s_j^{10}. \end{aligned} \quad (3.13)$$

Here $s_j = c_j/h_{j+1}$, c_j being the crack depth.

The requirements (3.7) form a linear algebraic system with the determinant Δ . Since this is a homogeneous system it has a non-trivial solution if and only if $\Delta = 0$. This equation enables to determine the natural frequencies for an arch containing a crack.

The results of calculations are presented in Fig.13—Fig.16 for the arch with $R = 1 \text{ m}$, $\beta = 1 \text{ rad}$, $h_0 = 0.02 \text{ m}$, $h_1 = 0.01 \text{ m}$.

The natural frequency versus α is depicted in Fig.13 for different extensions of the crack. The upper curve corresponds to the undamaged arch (with $s = 0$).

The dependence of natural frequency on the radius of the arch is shown in Fig.14 for different radii of the arch.

The natural frequency versus the step location is presented in Fig.15 and Fig.16 for the cases $s = 0$ and $s = 0.8$, respectively.

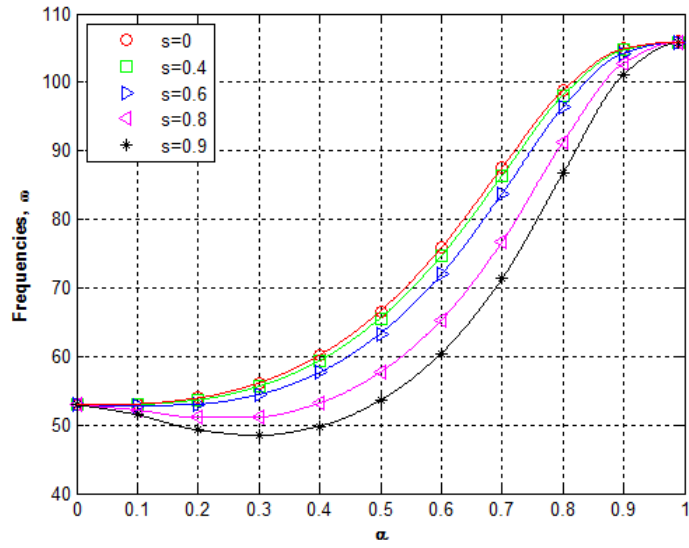


Figure 13: Natural frequency vs. depth of the crack

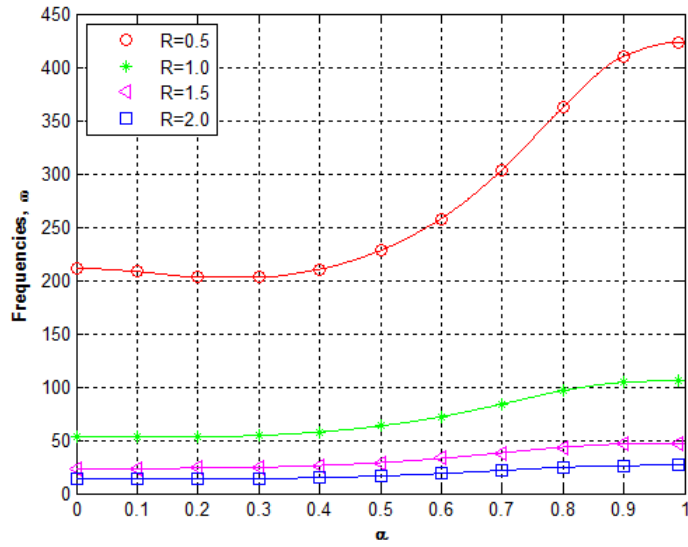


Figure 14: Natural frequency vs. radius of the arch

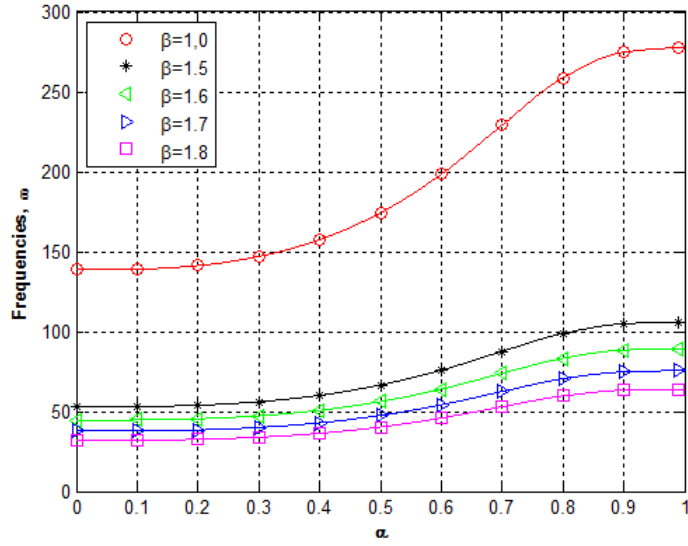


Figure 15: Natural frequency vs. step location ($s = 0$)

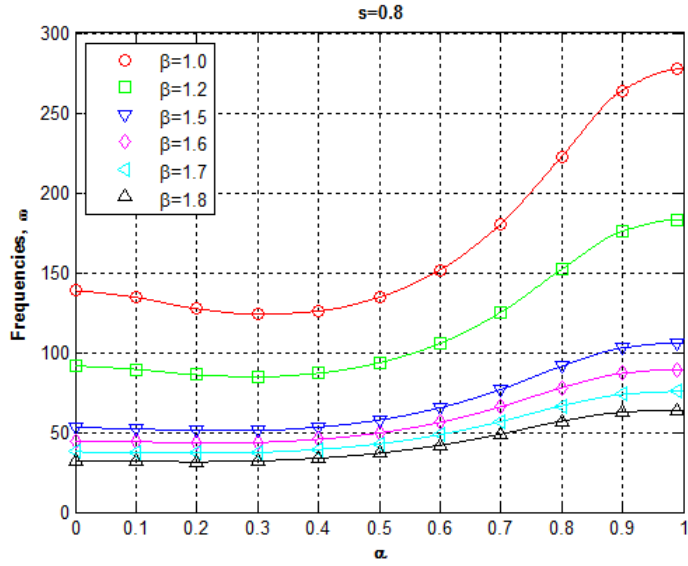


Figure 16: Natural frequency vs. step location ($s = 0.8$)

4 Free vibrations of tubular elastic arches

Let us consider the dynamic behavior of an elastic arch of rectangular uniform or non-uniform cross section. Here we will use capital letters to denote external dimensions of the cross-section of the arch and small letter for the internal dimensions. It is assumed that the external dimensions of the cross section are B (width) and H (height). Let b and h be for internal width and height, respectively.

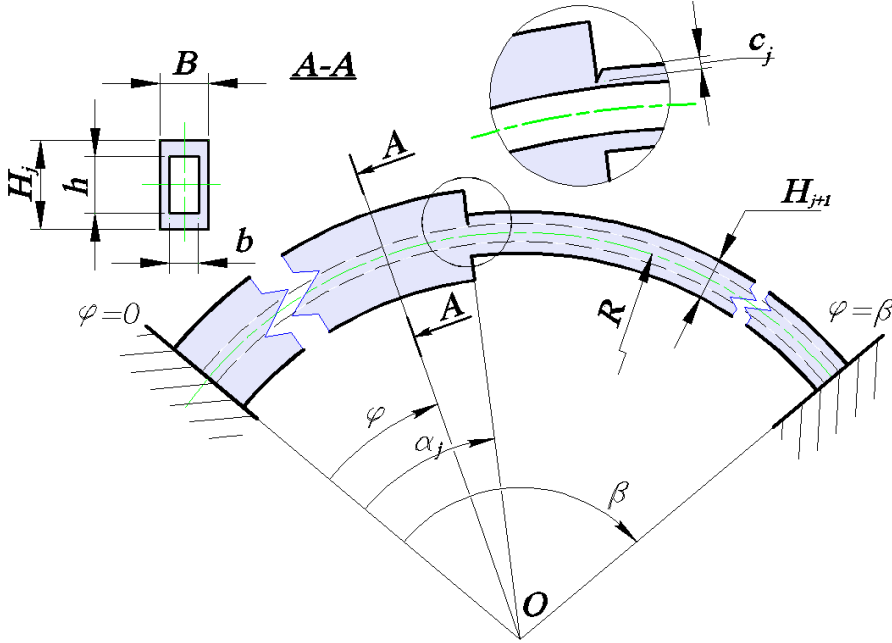


Figure 17: Stepped arch with tubular cross-sections

We assume that the height H is piece-wise constant, e.g.

$$H = H_j, \quad \varphi \in (\alpha_j, \alpha_{j+1}) \quad (4.1)$$

for $j = 0, \dots, n$. However, the radius of the arch $R = \text{const}$ and $B = \text{const}$. Let the arch be a circular arch with radius R and central angle β . (see Fig.17). For the sake of simplicity one can take $\alpha_0 = 0$, $\alpha_{n+1} = \beta$.

The cross sectional area of the arch is $S_j = BH_j$ for $j = 0, \dots, n$. However, the area of these parts which are occupied with the material is

$$S_j = BH_j - bh \quad (4.2)$$

for $\varphi \in (\alpha_j, \alpha_{j+1})$ $j = 0, \dots, n$.

The arch is weakened with cracks of length c_j locating at $\varphi = \alpha_j$, ($j = 1, \dots, n$).

The goal of the study is to define frequencies of natural vibrations and to analyze the sensitivity of eigenfrequencies on the crack parameters.

The basic equations used in the case remain valid.

The Hooke's law for a stepped arch can be presented as

$$M = -\frac{D_j}{R^2}(W + W'') \quad (4.3)$$

for $\varphi \in (\alpha_j, \alpha_{j+1})$ $j = 0, \dots, n$. Here

$$D_j = -E \iint_A z^2 dx dz, \quad (4.4)$$

where A is domain of integration and depends on the configuration of the cross section of the arch.

It appears that the system of equilibrium equations (2.1) can be converted in a single third order equation. According to the first equation in (2.1) and (2.6)

$$N' = -\frac{M'}{R}. \quad (4.5)$$

Differentiating the last equality in (2.1) with respect to φ and taking $Q = -N'$ and (4.5) into account one obtains

$$M''' + M' = \bar{\mu}_j R^2 \ddot{W}' \quad (4.6)$$

for $\varphi \in (\alpha_j, \alpha_{j+1})$ $j = 0, \dots, n$.

Substituting the bending moment from (4.3) into (4.6) one reaches to the equation

$$D_j(W^V + 2W''' + W') + \rho S_j R^4 \ddot{W}' = 0 \quad (4.7)$$

for $\varphi \in (\alpha_j, \alpha_{j+1})$ $j = 0, \dots, n$; where $\bar{\mu}_j = \rho S_j$, ρ being the density of the material.

The equation (4.7) can be solved by the method of separation of variables (see section 2.). Substituting (2.11) in (4.7) results in the equation

$$D_j(w^V + 2w''' + w') - w' \cdot \rho S_j R^4 \omega^2 = 0. \quad (4.8)$$

The latter can be presented as

$$w^V + 2w''' + w'(1 - k_j^2) = 0, \quad (4.9)$$

where

$$k_j = \sqrt{\frac{\rho S_j R^4 \omega^2}{D_j}}. \quad (4.10)$$

The characteristic equation of (4.9) has the form

$$\lambda_j^5 + 2\lambda_j^3 + \lambda_j(1 - k_j^2) = 0. \quad (4.11)$$

The roots of (4.11) are

$$\begin{aligned} \lambda_{j1} &= 0, \\ \lambda_{j2,3} &= \pm \sqrt{1 - k_j}, \\ \lambda_{j4,5} &= \pm i \sqrt{1 + k_j}, \end{aligned} \quad (4.12)$$

where i is the imaginary unit.

The general solution of (4.9) is

$$C_{1j} \cosh(\mu_j \varphi) + C_{2j} \sinh(\mu_j \varphi) + C_{3j} \cos(\nu_j \varphi) + C_{4j} \sin(\nu_j \varphi) + C_{5j}. \quad (4.13)$$

Here $\mu_j = \sqrt{1 - k_j}$, $\nu_j = \sqrt{1 + k_j}$.

The slope of the deflection is considered as a discontinuous quantity at the cross section with cracks.

In the case of hollow sectional beams instead of (2.22) one has

$$\sigma_j = \frac{6M_j H_j}{B_j H_j^3 - b h^3}. \quad (4.14)$$

The differential equation (here $s_j = c_j/H_j$) (2.22) takes the form

$$\frac{dC_j}{ds_j} = \frac{72H_j^4 B_j}{E'(B_j H_j^3 - b h^3)^2} s_j F^2(s_j). \quad (4.15)$$

The equation (4.15) with the initial condition $C_j(0) = 0$ has the solution

$$C_j = \frac{72H_j^4 B_j}{E'(B_j H_j^3 - b h^3)^2} f(s_j) \quad (4.16)$$

in the case where $c_j \leq \frac{1}{2}(H_j - h)$. Here the function $f(s_j)$ is defined by (3.13). Summarizing the results obtained above one can present the jump conditions for the slope of the deflection at $\varphi = \alpha_j$ as

$$\theta_j = -\frac{1}{R^2}C_j D_j(\alpha_j + 0) \left(w(\alpha_j + 0) + w''(\alpha_j + 0) \right), \quad (4.17)$$

where D_j and C_j are defined by (4.4) and (4.16), respectively. Note that (4.15), (4.16) apply for small cracks; in the case of larger cracks (4.15) and (4.16) must be modified suitably.

Some numerical results of calculation using the computer code MATLAB are presented in Fig. 18—Fig. 21.

The frequency of free vibrations is presented in Fig. 18 for various radii of the arch. The arch is made of a mild steel with $E = 2.1 \cdot 10^{11} \text{ Pa}$.

The natural frequency versus depth of the crack can be seen in Fig.19. The dimensions of the arch are $R = 1 \text{ m}$, $h_0 = 0.02 \text{ m}$, $h_1 = 0.01 \text{ m}$; material mild steel.

Natural frequencies of arches made of various materials are depicted in Fig. 20. Here $R = 1 \text{ m}$, $h_0 = 0.02 \text{ m}$, $h_1 = 0.01 \text{ m}$. Characteristics of the materials are shown in Table 4.1 below.

Table 4.1. Constants of different materials

	$E, (Pa)$	$\rho, (kg/m^3)$	ν
Steel	$2.1 \cdot 10^{11}$	7865	0.3
Aluminum	$7 \cdot 10^9$	2700	0.33
Glass	$7 \cdot 10^9$	2600	0.25
Concrete	$1.7 \cdot 10^9$	1500	0.15
Polystirol	$1.5 \cdot 10^9$	30	0.1

It can be seen that curves corresponding to aluminum, glass and steel are quite close to each other.

Different curves in Fig.21 correspond to different values of the ratio of thicknesses $\gamma = h_1/h_0$.

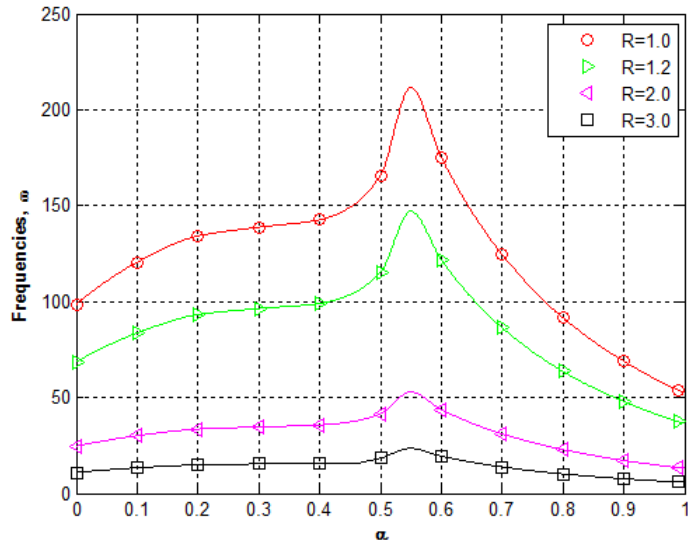


Figure 18: Natural frequency vs. radius of the arch

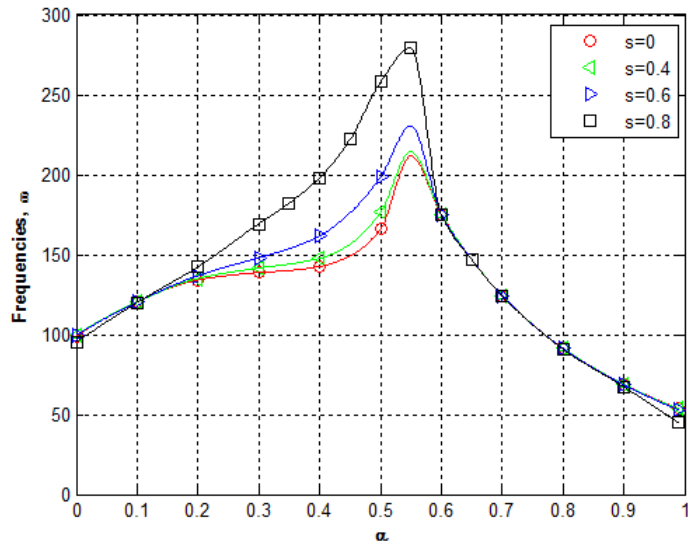


Figure 19: Natural frequency vs. depth of the crack

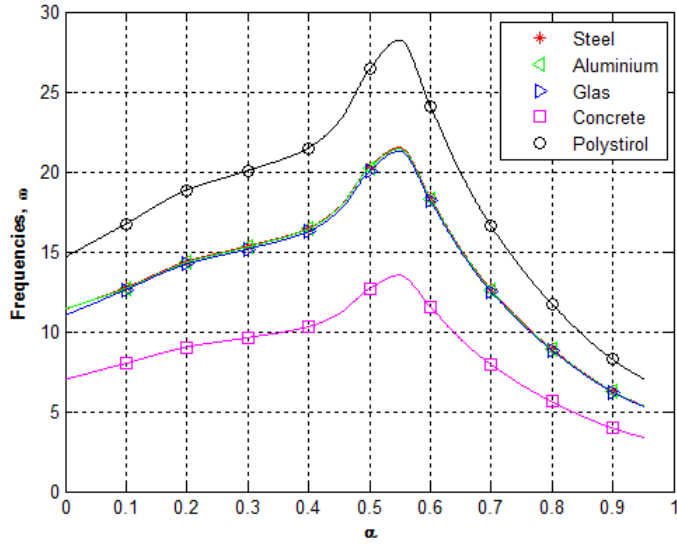


Figure 20: Natural frequency vs. material of the arch made

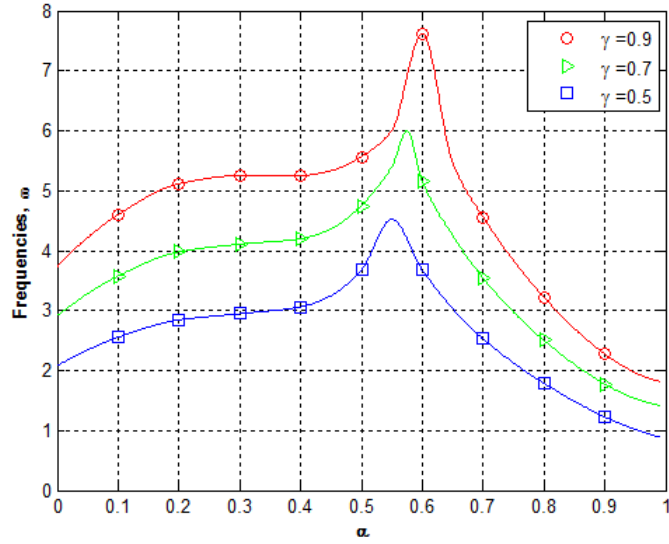


Figure 21: Natural frequency vs. ratio of the step $\gamma = \frac{h_1}{h_0}$

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Summary

In the present dissertation, the natural frequencies of free vibrations of circular elastic arches are studied. The circular elastic arches, in which each segment has a piece-wise constant thickness, are under consideration. It is assumed that stable crack-like surface defects occur at the re-entrant corners of the steps. The propagation of the cracks is neglected. The cracks are considered to be stationary surface cracks which have not fully penetrated the arch thickness. Combining the methods of the theory of elastic plates and shells and the theory of linear elastic fracture mechanics, a new method for determining the natural frequencies of elastic structures is developed in the thesis. The dissertation is based on the five articles of the author (four of which have been published). The dissertation consists of the review of the obtained results, the copies of the papers, the list of literature and the CV of the author. The review starts with the consideration of the background of the study of free vibrations, then overview of the used literature is provided and the aim of the thesis is presented (these parts form the Introduction), which is followed by the three main sections. In the second section of the dissertation, the method of finding the natural frequencies of arches with a constant thickness is developed. Two problems are discussed: the first one regards the arches made of homogeneous materials and containing surface cracks; the second case covers laminated arches without cracks. In the both cases, simply supported and clamped arches are studied. Because of the small amplitude of the oscillations in the case of free vibration, the material of the arches is assumed to be a purely elastic material and the hypothesis of Kirchhoff is considered applicable. A refined version of the classical bending theory is employed. The influence of cracks on the vibrational characteristics is taken into account with the use of the model of distributed line springs. The latter employs the stress intensity coefficient known in the elastic fracture mechanics. In the third section, stepped arches without and with cracks are studied. In the fourth section of the dissertation, the developed method is used for determination of the natural frequency of free vibration in case of hollow elastic arches. The stepped arches under consideration are assumed to be clamped at both ends. The influence of the geometrical parameters and material properties on the vibration of the arches is analysed. The method developed in the dissertation can be used in the non-destructive testing of structures.

Kokkuvõte

Pragudega elastsete astmeliste kaarte omavõnkumised

Käesolevas väitekirjas vaadeldakse elastsete astmeliste kaarte omavõnkumisi. Vaatluse all olevad kaared on tükiti konstantse paksusega ja konstantse laiusega. Kaarte tükiti konstantse paksuse muutumise kohtades asuvad praod. Eeldatakse, et praod on stabiilsed ja konstantse pikkusega. Kasutades elastsete plaatide ja koorikute teooria ning lineaarse purunemismehaanika meetodeid on välja töötatud analüütilis-numbriline meetod kaarte vabavõnkumiste sageduste määramiseks. Antud dissertatsioon põhineb autori viiel teaduslikul artiklil, millest neli on avaldatud ja üks ametlikult vastu võetud. Väitekirja koosneb kokkuvõtvast ülevaateartiklist, publitseeritud ja trükki suunatud artiklite koopiatest ja kokkuvõttest. Lisatud on autori elulookirjeldus.

Väitekirja sissejuhatavas osas on esitatud kirjanduse ülevaade ning on kirjeldatud väitekirja eesmärgid ja struktuuri. Töös uuritakse kaarte vabavõnkumisi konstantse ja tükiti konstantse paksuse korral. Esiteks vaadeldakse pinnalt lähtuva praoga homogeenest materjalist kaart, seejärel uuritakse lamineeritud elastse kaare vabavõnkumisi nii vabalt toetatud kui ka jäigalt kinnitatud kaare korral. Eeldatakse, et kehtivad Kirchhoffi hüpoteesid. Prao mõju vabavõnkumiste sagedusele modelleeritakse jaotatud lineaarse vedru mudeli abil. Vastavalt sellele on praio mõju kaare omavõnkumisele seotud lokaalse järeleandlikkuse koefitsiendi ning pingete intensiivsuse koefitsiendiga, mis arvutatakse purunemismehaanika meetoditega. Uuritakse ka astmeliste kaarte vabavõnkumisi. Vaadeldud on nii pragudega kaari kui ka ilma defektideta kaari. Dissertatsiooni viimane osa on pühendatud jäigalt kinnitatud õõnsustega kaarte vabavõnkumiste uurimisele. Kõikidel juhtudel on omavõnkumiste uurimiseks konstrueeritud analüütilis-numbrilised meetodid, mis põhinevad klassikalise plaatide ja koorikute teooria ning purunemismehaanika võrranditel ja kriteeriumitel.

Välja töötatud meetod võib leida kasutamist konstruktsioonide mittepurus-taval katsetamisel.

Acknowledgements

I wish to express my deepest gratitude to my supervisor Prof. Jaan Lellep for all the ideas, advice and continuous support he has generously provided throughout my doctoral study. I especially appreciate the personal wisdom and composure he showed in all these years of my studying.

My whole-hearted thanks go to my colleagues at the University of Tartu and the Estonian University of Life Sciences. Without their support and participation this work would not be possible.

I want to thank all the people, who have assisted me in this important period of my life the doctoral study. They have been many, and each of them has contributed to my professional formation.

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List of publications

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